

A Refinement of the Crystal Structure of Glycine*

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The crystal structure of glycine, originally determined by Albrecht & Corey in 1939, has been refined on the basis of complete three-dimensional intensity data obtained from molybdenum X-radiation. The structural parameters for the heavy atoms, including anisotropic temperature factors, were refined by difference maps and least-squares methods; the positions of the hydrogen atoms were located from a difference map. The final R factor was 0.063 for 1867 observed reflections and the standard deviations in the positional parameters of the heavy atoms are about 0.001 Å.

The bond distances and angles are close to those reported by Albrecht & Corey, with the exception of the C-N bond, which is found to be 1.474 Å rather than 1.39 Å. The apparent C-H and N-H distances are about 0.9 Å.

Introduction

The crystal structure of glycine, $\text{H}_3^+\text{NCH}_2\text{COO}^-$, was first determined by Albrecht & Corey (1939). The structure was solved by Patterson and Patterson-Harker maps, and was refined by trial-and-error methods. Although three-dimensional intensity data (copper radiation) were used for the Patterson maps, the refinement was based entirely on zonal data. Unfortunately, for a small molecule glycine is surprisingly ill-suited for two-dimensional refinement, since there is no direction of projection which does not involve serious overlap.

A surprising result obtained by Albrecht & Corey was the value 1.39 ± 0.02 Å for the C-N distance—a value shorter by about 0.08 Å than that normally found in amino acids and peptides. Furthermore, their results indicated that, although two of the hydrogen atoms of the NH_3 group are involved in normal hydrogen bonding with neighboring oxygen atoms, the third hydrogen atom forms a 'bifurcated' bond with two adjacent oxygen atoms. In view of these surprising features, a redetermination of the structure of this simplest of all amino acids was in order.

It has seemed worth while to put some special effort into this redetermination, with a view not only towards obtaining accurate parameters for the heavier atoms and locating, if possible, the hydrogen atoms, but also towards learning something concerning the accuracy to which ordinary X-ray diffraction techniques may be pushed.

Experimental

(i) Refinement of the unit-cell parameters

Large single crystals of glycine (the usual, or α , form) were readily grown by slow evaporation of

aqueous solutions. They were elongated in the c direction and showed pronounced cleavage perpendicular to b . One crystal was ground to a cylinder about 3 mm. in length and 1.2 mm. in diameter; this was used for photographs about the c axis. A second crystal was shaped into an approximate sphere about 1 mm. in diameter by rolling it between thumb and forefinger, the slight moisture on the fingers being sufficient to dissolve the rough edges; this crystal was used for photographs about the a , b , and $[101]$ axes.

Accurate values for the unit-cell dimensions were obtained from three Straumanis-type rotation photographs taken about the b , c , and $[101]$ axes, respectively. The copper X-ray tube was badly contaminated with iron, so that a number of the reflections used in the determination of cell constants were due to iron rather than copper radiation. The photographs were measured with the help of a traveling microscope and a calibrated steel scale; the effective camera radius varied from 4.981 to 5.002 cm. for the three photographs. No eccentricity or absorption effect was apparent.

Two or three high-angle reflections were chosen on each film and were indexed on the basis of the unit-cell dimensions reported by Albrecht & Corey. In all, seven reflections having Bragg angles greater than 70° were chosen; of these, three were of copper radiation and

Table 1. *Spacing data for glycine*

Rotation axis	hkl	Radiation	$(\sin^2 \theta)_o$	$(\sin^2 \theta)_c$	Weight
c	{	470 Fe $K\alpha_1$	0.98760	0.98759	81
		1,12,0 Fe $K\alpha_1$	0.98330	0.98326	38*
		1,15,0 Cu $K\alpha_1$	0.95783	0.95794	24
$[101]$	{	2,14, $\bar{2}$ Cu $K\alpha_1$	0.93646	0.93636	17
		4,10, $\bar{4}$ Cu $K\alpha_1$	0.91352	0.91356	13
b	{	204 Fe $K\alpha_1$	0.98058	0.98058	52
		503 Fe $K\alpha_1$	0.93797	0.93797	17

* This reflection was very weak and the reported weight includes a factor of 0.64 as well as the standard factor $1/\sin^2 2\theta$.

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four were of iron. The spacings of the $K\alpha_1$ components were measured and used in a least-squares determination of the four unit-cell parameters. Values of $\sin^2 \theta$ for these seven reflections and their assigned weights are listed in Table 1; the output cell parameters are

Table 2. *Unit-cell dimensions for glycine*
(λ (Cu $K\alpha_1$) = 1.54051 Å; λ (Fe $K\alpha_1$) = 1.93597 Å)

	This investigation	Albrecht & Corey
<i>a</i>	5.1020 ± 0.0008* Å	5.10 Å
<i>b</i>	11.9709 ± 0.0017* Å	11.96 Å
<i>c</i>	5.4575 ± 0.0015* Å	5.45 Å
β	111° 42.3' ± 1.0*	111° 38'

* Estimated limits of error set at about five times the standard deviations calculated in the least-squares treatment.

listed in Table 2, together with the values reported by Albrecht & Corey.

The density calculated on the basis of four molecules in the unit cell is 1.610 g.cm.⁻³; the observed density has been reported in the literature as 1.606 (Low & Richards, 1952) and as 1.607 (Curtius, 1882) g.cm.⁻³.

(ii) Collection of intensity data

Equi-inclination Weissenberg photographs were taken with Mo $K\alpha$ radiation for layer lines zero through seven from crystals rotated about *c*, *a*, and [101]. Three films were used for each exposure, with thin copper foil interspersed to bring the film factor up to about 3.5; for most layer lines, two exposures of different duration were taken. Intensities were estimated by visual comparison with intensity strips prepared from the same crystals. Empirical film factors were obtained, and found to vary in the expected manner with the equi-inclination angle.

The intensities were corrected for Lorentz and polarization effects. Correlation factors for the various

exposures were then obtained by comparing intensity values from each set of films with those from all sets from the other two rotation axes and taking subjectively-weighted averages. Finally, F^2 values, on an arbitrary scale, were obtained by taking averages (again, subjectively weighted) of the values observed about the three different axes. All reflections within the molybdenum sphere—about 5400—were covered, of which 1919 were strong enough to be observed. Of these, 965 were observed on photographs taken about all three axes, 588 were observed about two axes, and 366 were observed about only one axis. Values for the observed intensity ranged from about 2 to 11,000.

In order to assess the uncertainties in the observed structure factors, comparisons were made among the F^2 values for all reflections which were observed about more than one axis. The estimated relative standard deviations averaged about 5% for reflections of moderate intensity, increasing to about 10% for the strong reflections and to about 20% for the very weak reflections. This pattern of observational discrepancies was the basis of the weighting function used in the least-squares refinement of the positional parameters.

Refinement of the atomic parameters

The atomic positional parameters were refined by means of electron-density and difference projections on to (010), (100), and (101); through six three-dimensional structure-factor and difference-map cycles; and, finally, through two structure-factor least-squares cycles. Individual anisotropic temperature factors were included for the heavy atoms, beginning with the two-dimensional calculations; they were not optimized in the final least-squares treatments. Positional parameters for the hydrogen atoms were obtained from a three-dimensional difference map.

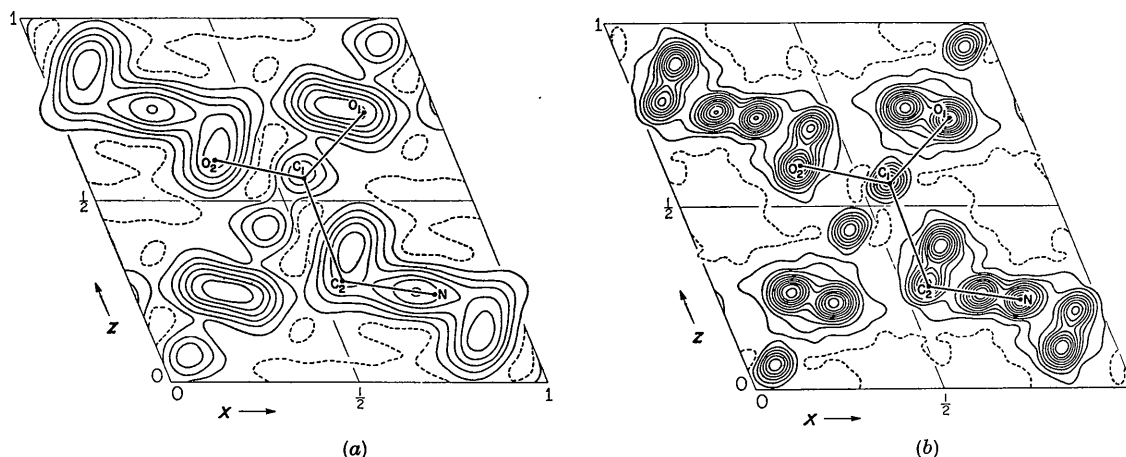


Fig. 1. Electron-density projections on to (010): (a) calculated from the copper intensity data of Albrecht & Corey; (b) calculated from the molybdenum data of the present investigation. The contours are at intervals of 2 e.Å⁻²; the 2 e.Å⁻² contour is broken.

Table 3. *Atomic parameters obtained from two-dimensional refinements*

	x	y	z	α	β	γ	δ	ε	η
C ₁	0.0760	0.1247	0.0642	0.0125	0.0025	0.0100	0	0.0125	0
C ₂	0.0655	0.1453	0.7880	0.0125	0.0030	0.0105	0	0.0130	0
N	0.3030	0.0905	0.7410	0.0120	0.0035	0.0075	0	0.0075	0
O ₁	0.3060	0.0936	0.2350	0.0145	0.0040	0.0115	0	0.0120	0
O ₂	0.8538	0.1400	0.1082	0.0175	0.0045	0.0150	0	0.0175	0

(i) *Two-dimensional refinements*

The starting parameters for the two-dimensional refinements were those of Albrecht & Corey (1939). From these parameters, signs for the structure factors of the type $h0l$ were calculated; these signs were used in the preparation of an electron-density projection on to (010), in which all observed $h0l$ terms were included. This projection is shown in Fig. 1(b). For comparison, a second electron-density projection on to (010), based on the structure factors observed by Albrecht & Corey and listed in their Table 2, was calculated; it is shown in Fig. 1(a). The increased resolution obtained from the molybdenum intensity data of the present investigation, as compared with the copper data of Albrecht & Corey, is clearly apparent.

The x and z parameters indicated by the (010) projection were used in the calculation of $h0l$ structure factors, and approximate scale and isotropic temperature factors were assigned. The R factor for these 69 reflections was 0.149. A subsequent (010) difference projection showed only small shifts for the positional parameters, but gave indications of highly anisotropic temperature factors, particularly for the oxygen atoms. Accordingly, preliminary values for the individual temperature-factor anisotropies were assigned, the temperature factors being expressed in the form

$$T_i = \exp -(\alpha_i h^2 + \beta_i k^2 + \gamma_i l^2 + \delta_i hk + \varepsilon_i hl + \eta_i kl).$$

Attention was then turned to the $0kl$ reflections, from which an electron-density projection on to (100) and two structure-factor difference-map cycles were calculated. A second difference projection on to (010) and a structure-factor difference-map cycle, using data of the type hkh , completed the two-dimensional work. The atomic parameters at this stage of refinement are listed in Table 3.

(ii) *Three-dimensional refinements*

With the exception of the final structure-factor least-squares cycles, which were carried out on a Datatron computer, the three-dimensional calculations were made on conventional IBM equipment, including a 604 Electronic Calculating unit. The first set of structure factors, calculated from the parameters in Table 3, led to an R factor of 0.131 for the observed reflections; they were used in the calculation of the first three-dimensional difference map. This and the subsequent difference maps were evaluated at inter-

vals of $a/30$, $b/60$, and $c/30$, by use of the M -card system; accordingly, it was necessary to omit about 100 observed reflections having h or l greater than 10 or k greater than 20.

The first difference map indicated shifts in the positional parameters averaging about 0.003 Å; more significant were the indicated changes in the individual temperature factors. In particular, it was evident that the temperature-factor anisotropies should be markedly increased for the oxygen atoms. Approximate values for the changes in the scale and temperature-factor parameters were obtained by the method of Leung, Marsh & Schomaker (1957)*, and a second set of structure factors led to an R factor of 0.117. Another difference-map structure-factor cycle reduced R to 0.109.

At this point, the contributions of the hydrogen atoms were calculated. The parameters chosen for the two methylene hydrogen atoms were calculated on the basis of tetrahedral bond angles and C-H distances of 1.00 Å; the parameters for the ammonium hydrogen atoms were those indicated by the first three-dimensional difference map, which led to N-H distances of about 0.88 Å and to a roughly tetrahedral configuration about the nitrogen atom. An isotropic temperature factor with $B = 2.0$ Å² was assigned to each of the five hydrogen atoms. The inclusion of these hydrogen-atom contributions led to a significant improvement in the agreement of the low-order data.

It was now apparent that the observed structure factors for a number of strong reflections were too small, indicating extinction effects. Accordingly, the F_o values were multiplied by the correction factor $\sqrt{\{1/(1-KI_o)\}}$, where I_o is the observed intensity; the value for K was determined empirically. This correction was applied to 119 observed structure factors, of which 63 were changed by more than 1% and 8 by more than 10%. After including these extinction corrections and the hydrogen-atom contributions, the R factor was 0.092.

A third difference map was then calculated; it was much flatter than the first two, but again indicated that the major source of the discrepancies between observed and calculated structure factors lay in the assignment of temperature-factor parameters. These were re-adjusted and the structure factors were recalculated; R dropped to 0.077. A fourth difference map indicated only small changes in the temperature-

* The equations for the temperature-factor corrections are not exact when the input atoms are not isotropic.

factor and positional parameters; the next set of structure factors led to an R factor of 0.072.

A difference map was then calculated with F_c values from which the hydrogen-atom contributions had been omitted, and including only those terms with h and l less than 5 and k less than 10. The positions of the five hydrogen atoms were clearly apparent on this map; they lay in regions with maximum electron density ranging from 0.75 to 0.90 e.Å⁻³, whereas there were no other regions of density greater than 0.3 e.Å⁻³. A composite drawing of this map is shown

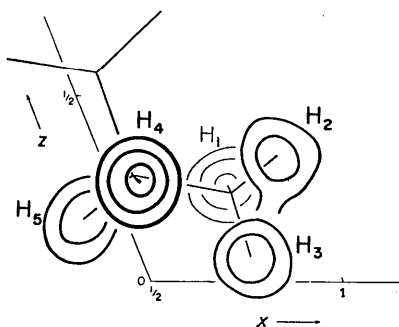


Fig. 2. A composite drawing of the three-dimensional difference map in which the hydrogen-atom contributions were omitted from the F_c values. The contours are at intervals of $\frac{1}{4}$ e.Å⁻³, beginning with the $\frac{1}{4}$ e.Å⁻³ contour.

in Fig. 2, and the indicated hydrogen-atom parameters are listed in Table 4.

Table 4. *The final hydrogen-atom parameters*

Atom	x	y	z
H ₁	0.286	0.100	0.570
H ₂	0.457	0.119	0.837
H ₃	0.298	0.020	0.763
H ₄	0.080	0.220	0.771
H ₅	0.898	0.118	0.671

At this stage of the investigation, a Datatron digital computer became available. The computer had previously been programmed for structure-factor and least-squares calculations on structures with point groups 2 or m (Pasternak, 1956), and the program was readily modified to accommodate the remaining monoclinic space groups. This program accepts complete anisotropic temperature factors for each atom, but only the positional parameters are optimized.

The weighting function used in the least-squares calculations was derived from the estimated standard

deviations in the F_o values (see section 'Experimental'), the values for \sqrt{w} being taken inversely proportional to the expected uncertainty in F_o . The explicit weighting function was:

$$\begin{aligned} \sqrt{w} &= 35/F_o(1+0.0032F_o^2) & \text{for } F_o \geq 5.3, \\ \sqrt{w} &= 6.0 & \text{for } F_o \leq 5.3. \end{aligned}$$

Two additional factors were included in \sqrt{w} : the factor \sqrt{n} , where n ($= 1, 2,$ or 3) is the number of rotation axes about which the observation was made, and a subjective factor (sometimes equal to zero) for observations which, for reasons of extinction, spot shape, contamination with white-radiation streaks, or extreme weakness, were of doubtful validity. Finally, for computational convenience, the values of \sqrt{w} were rounded off to a single integer.

The first set of structure factors calculated on the Datatron computer led to an R factor—including only those reflections with $\sqrt{w} \neq 0$ —of 0.076. The structure factors differed slightly from those of the previous set for three reasons: (1) the hydrogen-atom parameters had been changed slightly; (2) the atomic form factors of Hoerni & Ibers (1954) were used in place of those of McWeeny; (3) the atomic form factors were stored in the computer as linear interpolation tables, and for a few very small regions in reciprocal space the interpolated form factors could be in error by nearly 5%. Because of a programming error, the indicated least-squares shifts from this calculation were not sensible; accordingly, the structure factors were used in the calculation of a final three-dimensional difference map. This map indicated shifts in the positional parameters averaging about 0.001 Å and shifts in the temperature-factor parameters averaging about 2%; furthermore, it indicated that the change to Hoerni-Ibers form factors had somewhat improved the over-all pattern of agreement. The most significant feature of this map—and one which was also apparent on the previous maps—was the presence of a negative region of about 0.6 e.Å⁻³ at the position of the nitrogen atom and of two smaller positive regions at the positions of the oxygen atoms, indicating that the use of form factors for neutral N and O atoms was inappropriate. The implied charge effect was not taken into account in the final structure-factor calculations, although it would apparently account for the relatively poor agreement in some low-order reflections.

The improved temperature-factor parameters were

Table 5. *Final positional and temperature-factor parameters*

(The temperature factor for each atom is of the form

$$T_i = \exp -(\alpha_i h^2 + \beta_i k^2 + \gamma_i l^2 + \delta_i hk + \epsilon_i hl + \eta_i kl)$$

Atom	x	y	z	$10^4\alpha$	$10^4\beta$	$10^4\gamma$	$10^4\delta$	$10^4\epsilon$	$10^4\eta$
C ₁	0.07542	0.12478	0.06605	157	22	115	-10	150	-10
C ₂	0.06536	0.14499	0.78711	156	32	118	20	130	16
N	0.30135	0.08980	0.74113	164	38	108	34	130	12
O ₁	0.30583	0.09427	0.23553	205	50	115	45	138	28
O ₂	0.85224	0.14154	0.10711	185	56	190	8	230	-20

Table 6 (cont.)

2,k,10	3,k,3	5,k,2	4 72 86	12 108 66	10 876 608	5,k,7	5,k,1	6 481 -61
0 216 -219	0 495 -554	1 2949 -2963	5 819 812	12 297 294	11 279 304	1 72 44	0 288 -307	7 7 -739
1 284 307	1 351 336	2 179 1766	6 108 112	15 171 162	12 261 -270	2 1143 -1135	1 240 -555	8 171 -201
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4 232 214	4 459 495	5 1045 1045	10 <108 -60	24 <135 -26	16 171 163	5 99 -93	4 279 303	12 <108 79
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	72 <135 -51	73 135 -51	74 135 -51	228 <135 63	21 153 -111	73 126 -108	72 126 -108	80 126 108
	73 <135 -51	74 135 -51	75 135 -51	231 <135 63	21 153 -111	74 126 -108	73 126 -108	81 126 108
	74 <135 -51	75 135 -51	76 135 -51	234 <135 63	21 153 -111	75 126 -108	74 126 -108	82 126 108
	75 <135 -51	76 135 -51	77 135 -51	237 <135 63	21 153 -111	76 126 -108	75 126 -108	83 126 108
	76 <135 -51	77 135 -51	78 135 -51	240 <135 63	21 153 -111	77 126 -108	76 126 -108	84 126 108
	77 <135 -51	78 135 -51	79 135 -51	243 <135 63	21 153 -111	78 126 -108	77 126 -108	85 126 108
	78 <135 -51	79 135 -51	80 135 -51	246 <135 63	21 153 -111	79 126 -108	78 126 -108	86 126 108
	79 <135 -51	80 135 -51	81 135 -51	249 <135 63	21 153 -111	80 126 -108	79 126 -108	87 126 108
	80 <135 -51	81 135 -51	82 135 -51	252 <135 63	21 153 -111	81 126 -108	80 126 -108	88 126 108
	81 <135 -51	82 135 -51	83 135 -51	255 <135 63	21 153 -111	82 126 -108	81 126 -108	89 126 108
	82 <135 -51	83 135 -51	84 135 -51	258 <135 63	21 153 -111	83 126 -108	82 126 -108	90 126 108
	83 <135 -51	84 135 -51	85 135 -51	261 <135 63	21 153 -111	84 126 -108	83 126 -108	91 126 108
	84 <135 -51	85 135 -51	86 135 -51	264 <135 63	21 153 -111	85 126 -108	84 126 -108	92 126 108

A REFINEMENT OF THE CRYSTAL STRUCTURE OF GLYCINE

Table 6 (cont.)

5,k,7	3 351 357 4 481 44 5 630 638 6 268 267 7 268 267 8 153 157 9 153 157 10 153 157 11 153 157 12 153 157 13 153 157 14 153 157 15 153 157	5,k,5	1 261 255 2 413 98 3 486 98 4 198 170 5 153 158 6 153 158 7 477 482 8 283 281 9 117 128 10 180 184 11 288 301 12 171 170 13 144 137 14 180 182 15 180 182	5,k,9	1 225 246 2 134 139 3 144 157 4 189 191 5 207 207 6 4135 41 7 4135 68 8 282 267 9 135 121 10 415 41 11 415 41 12 216 222 13 153 35 14 162 136 15 415 65 16 134 127 17 415 14 18 135 122	7,k,7	0 540 556 1 180 192 2 4135 47 3 4135 68 4 489 489 5 288 287 6 126 138 7 4171 41 8 333 335 9 282 271 10 180 173 11 4135 36 12 180 182 13 117 114 14 180 141	8,k,2	0 171 194 1 207 218 2 171 145 3 4135 41 4 4135 41 5 4135 41 6 279 278 7 279 278 8 216 212 9 4153 3 10 198 194 11 4171 133 12 333 342 13 159 149 14 4189 16 15 153 141 16 279 278 17 4189 45 18 4189 60 19 4189 60 20 153 142	5<198 166 7 4198 36 8 297 294 9 4198 142 10 333 334 11 4198 30 12 315 317 13 4198 103 14 4198 116 15 4198 50 16 207 186	5,k,5	5 162 142 9 144 128 1 4153 109 2 360 399 3 207 258 4 225 258 5 4153 60 6 189 175 7 225 265 8 234 282 9 144 153 10 4153 110 11 180 204 12 162 139 13 153 147																																															
5,k,8	1 279 267 2 4135 42 3 4135 42 4 324 345 5 297 300 6 171 190 7 396 396 8 378 396 9 4135 45 10 4135 45 11 378 392 12 243 237 13 135 117 14 4153 26 15 180 185	6,k,1	1 72 117 2 542 544 3 481 37 4 216 236 5 144 153 6 342 355 7 4108 12 8 283 286 9 106 150 10 4135 45 11 104 120 12 243 237 13 135 117 14 144 144 15 180 185	5,k,4	0 486 543 1 126 124 2 234 236 3 378 375 4 216 236 5 144 153 6 342 355 7 4108 12 8 283 286 9 106 150 10 4135 45 11 104 120 12 243 237 13 135 117 14 144 144 15 180 185	5,k,10	0 384 317 1 4153 35 2 153 134 3 4153 10 4 171 179 5 126 148 6 261 262 7 4153 35 8 4153 35 9 207 198 10 180 186 11 135 138 12 198 201 13 160 145 14 4153 20 15 126 133 16 182 156	7,k,8	0 171 185 1 297 299 2 481 19 3 117 16 4 180 187 5 387 404 6 117 116 7 90 72 8 198 196 9 297 315 10 4135 15 11 171 128 12 180 159 13 162 136	8,k,3	1 189 185 2 315 314 3 189 189 4 4135 47 5 4171 105 6 279 285 7 4171 109 8 4153 119 9 180 180 10 171 123 11 207 195 12 171 123 13 171 95	5,k,7	0 315 323 1 99 147 2 4135 107 3 126 158 4 297 283 5 4135 94 6 198 234 7 4135 110 8 198 196 9 4135 49 10 180 226																																														
5,k,9	0 4135 86 1 135 129 2 4135 11 3 4135 74 4 4153 53 5 279 277 6 4135 23 7 198 194 8 4153 17 9 225 244 10 4135 15 11 198 200 12 4153 20 13 135 122 14 4153 7 15 171 164 16 4153 30 17 4153 42 18 126 118	6,k,2	0 270 241 1 288 211 2 4108 75 3 441 456 4 283 296 5 4108 76 6 189 196 7 216 221 8 207 202 9 270 294 10 252 256 11 198 194 12 126 138 13 243 226 14 150 159 15 225 218	5,k,5	1 198 195 2 99 24 3 477 468 4 108 21 5 144 144 6 153 144 7 477 481 8 207 202 9 270 294 10 252 256 11 198 194 12 126 138 13 243 226 14 150 159 15 225 218 16 180 189	6,k,3	1 162 181 2 279 112 3 252 237 4 4135 47 5 4153 19 6 207 202 7 243 231 8 4135 96 9 4153 89 10 171 170 11 153 178	5,k,10	1 135 151 2 297 302 3 4153 88 4 126 129 5 225 220 6 4153 37 7 279 257 8 135 110 9 4153 134 10 171 198 11 234 250 12 4153 4 13 4153 4 14 243 242 15 4153 49 16 165 165 17 4153 37 18 162 165	6,k,4	0 261 270 1 153 169 2 4135 70 3 414 478 4 180 212 5 4135 57 6 153 193 7 261 308 8 4171 99 9 4171 96 10 216 239 11 4135 49 12 4153 30 13 4171 78 14 189 195	5,k,6	0 792 767 1 153 149 2 108 204 3 486 440 4 486 491 5 261 255 6 4153 137 7 261 293 8 4108 27 9 360 367 10 4153 3 11 153 153 12 360 345 13 126 179 14 4135 14 15 4153 23 16 324 338 17 4153 29 18 180 212 19 4135 57 20 171 199	7,k,0	1 126 152 2 395 301 3 153 157 4 89 212 5 4171 167 6 4171 121 7 414 381 8 270 245 9 4153 40 10 4153 71 11 315 337 12 162 126 13 152 174 14 153 143 15 153 139	7,k,1	0 4153 8 1 153 152 2 4108 31 3 4108 4 4 180 17 5 234 227 6 4135 75 7 153 146 8 135 142 9 135 142 10 171 182 11 126 146	7,k,2	1 180 166 2 180 191 3 4171 102 4 198 214 5 24 269 6 171 132 7 279 325 8 279 297	7,k,3	0 216 155 1 4153 61 2 4135 23 3 481 49 4 189 184 5 117 128 6 99 65 7 126 168 8 171 167 9 171 171 10 136 159 11 144 144 12 153 152 13 144 117	7,k,4	0 384 280 1 4108 25 2 4108 33 3 423 421 4 279 284 5 72 127 6 180 185 7 558 591 8 279 290 9 198 175 10 198 175 11 369 326	7,k,5	0 468 502 1 353 352 2 135 146 3 378 393 4 387 397 5 252 253 6 333 382 7 279 298 8 153 152 9 171 166 10 242 250 11 126 157 12 4135 62 13 153 144 14 225 260	7,k,6	1 153 137 2 135 230 3 4108 24 4 324 347 5 4108 46 6 108 84 7 126 128 8 342 348 9 135 146 10 135 147 11 171 140 12 153 143 13 198 236 14 108 135 15 4153 64 16 4153 39 17 126 126 18 198 227	7,k,7	1 495 468 2 765 750 3 450 442 4 135 151 5 333 326 6 350 353 7 315 359 8 126 116 9 4153 109 10 4135 464 11 216 217 12 4153 36 13 4171 114 14 189 207 15 180 163	7,k,8	1 495 468 2 765 750 3 450 442 4 135 151 5 333 326 6 350 353 7 315 359 8 126 116 9 4153 109 10 4135 464 11 216 217 12 4153 36 13 4171 114 14 189 207 15 180 163	7,k,9	0 270 296 1 4108 108 2 4198 21 3 4198 82 4 297 279 5 4198 37 6 4198 40 7 4198 35 8 282 230 9 4198 281 10 4198 131 6 114 145 7 153 160 8 234 245 9 261 277 10 126 127 11 4153 32 12 4153 32 13 216 212 14 153 152 15 4153 70 16 4153 127 17 153 153 6 333 374	8,k,4	0 909 1026 1 480 473 2 4153 407 3 197 197 4 693 768 5 369 381 6 114 145 7 153 160 8 234 245 9 261 277 10 126 127 11 4153 32 12 4153 32 13 216 212 14 153 152 15 4153 70 16 4153 127 17 153 153 6 333 374	8,k,5	1 252 267 2 378 409 3 292 273 4 4153 34 5 144 132 6 351 353 7 171 181 8 1171 83 9 279 291 10 4153 34 11 297 281 12 4153 35 13 126 102 14 180 207	7,k,11	0 136 125 1 306 344 2 4153 113 3 279 291 4 4153 34 5 279 281 6 153 156 7 171 171	8,k,6	0 639 706 1 180 198 2 135 137 3 171 171 4 489 508 5 225 227 6 294 269 7 171 190 8 114 128 9 171 181 10 135 140 11 207 198 12 4153 103 13 135 85 14 153 116 15 171 152	8,k,7	0 639 706 1 180 198 2 135 137 3 171 171 4 489 508 5 225 227 6 294 269 7 171 190 8 114 128 9 171 181 10 135 140 11 207 198 12 4153 103 13 135 85 14 153 116 15 171 152	8,k,8	0 567 589 1 414 472 2 306 327 3 441 506 4 4108 34 5 252 276 6 180 223 7 369 399 8 4135 57 9 135 95 10 4135 99 11 4135 99 12 4135 61 13 135 85 14 153 116 15 171 152	8,k,9	0 297 293 1 297 318 2 4171 7 3 198 212 4 252 270 5 279 10 6 4171 22 7 4171 41 8 171 175 9 207 211 10 4153 26 11 126 78 12 4153 73 13 153 120	8,k,10	0 297 293 1 297 318 2 4171 7 3 198 212 4 252 270 5 279 10 6 4171 22 7 4171 41 8 171 175 9 207 211 10 4153 26 11 126 78 12 4153 73 13 153 120	8,k,11	1 171 132 2 270 298 3 171 160 4 180 188 5 180 140 6 180 79 7 207 149 8 189 190 9 159 164 10 144 123 11 162 81 12 4162 61 13 162 149 14 162 164	8,k,12	0 4153 70 1 162 135 2 4171 75 3 162 144 4 4153 2 5 126 58 6 198 196 7 4153 103 8 153 138 9 297 293 10 297 318 11 4171 7 12 198 212 13 252 270 14 279 10 15 4171 22 16 4171 41 17 171 175 18 207 211 19 4153 26 20 126 78 21 4153 73 22 153 120	8,k,13	1 216 262 2 153 128 3 100 221 4 4153 62 5 180 190	8,k,14	1 153 168 2 4153 78 3 216 220 4 315 356 5 4153 54 6 198 208 7 198 212 8 387 416 9 4153 40 10 252 269 11 189 173 12 207 233 13 153 54 14 207 213 15 153 124	8,k,15	0 4180 102 1 180 187 2 4180 22

used in two structure-factor least-squares cycles. After the first cycle the author exercised his privilege of checking the observed intensities of reflections showing poor agreement; a few of these observed intensities (or their indices) were found to be in error. The final structure-factor least-square calculation was then made; the maximum parameter shift was 0.0014 Å and the average shift was 0.0004 Å. The final parameters, which include the shifts indicated in the last least-squares calculation,* are listed in Table 5.

The observed and calculated structure factors are listed in Table 6.

Accuracy of the results

The final R factor, calculated for 1867 observed reflections of non-zero weight, is 0.063. Aside from the observational uncertainties, the major sources of disagreement between observed and calculated structure factors probably are, in order of importance: (1) residual-charge effects on the nitrogen and oxygen atoms; (2) incomplete refinement of the temperature-factor parameters, including those of the hydrogen atoms; (3) use of inappropriate form factors, due to bond effects (which could be partially, but not entirely, compensated for by the temperature-factor parameters) and to the linear interpolation scheme used in the Datatron program.

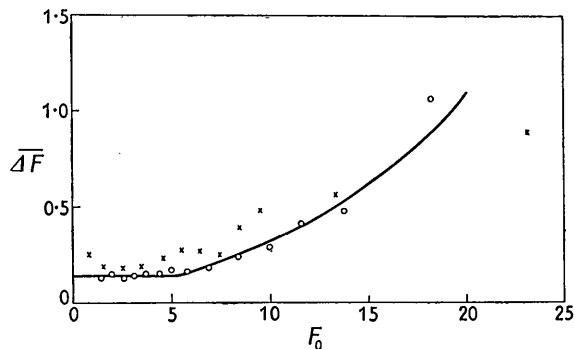


Fig. 3. Structure-factor discrepancies.

Discrepancies ΔF , averaged over all reflections within a range of F_0 values, are plotted against the mean F_0 value for that range. The crosses refer to the final structure-factor calculations; here, $\Delta F = |F_0 - F_c|$. The circles refer to the observational discrepancies implied by the disagreements between F^2 values measured about three different crystal axes i ; here, $\Delta F = \frac{|F_i^2 - \hat{F}^2|}{2F_0}$, where \hat{F}^2 ($\equiv F_0^2$) is the average of the three measurements F_i^2 . The solid line is proportional to the inverse of the weighting function $1/w$ used in the least-squares refinements. The units are electrons per unit cell.

* Only the diagonal terms of the normal equations are calculated in the least-squares routine, although in the present case the x - z cross terms would be large and, in view of the large temperature-factor anisotropies, not easily predictable. In the first least-squares, the magnitude of the x - z cross term was estimated; the observed effects of the parameter shifts on the least-squares totals were then used to calculate more accurate values for the cross terms which were included in the normal equations of the second least-squares.

A comparison of the final structure-factor agreement with that expected from the statistics of the observations is shown in Fig. 3. Here the discrepancy $|\Delta F|$ between F_0 and F_c , averaged over a group of reflections, is plotted against F_0 ; for comparison, the discrepancies arising from the statistics of the observations (see section 'Experimental') are also plotted. In addition, the pattern of uncertainties used as the basis of the weighting function is shown as a solid line in Fig. 3; this curve is proportional to the inverse of the weighting function $1/w$. It is apparent that the refinement is nearly, but not exactly, as complete as could be expected from the accuracies of the observations. Accordingly, the weighting scheme which was used in the least-squares refinement of the positional parameters, while not perfect, is clearly appropriate.

The standard deviations in the positional parameters of the carbon, nitrogen, and oxygen atoms, calculated from the residuals of the last structure-factor least-squares cycle, range from 0.0007 Å to 0.0010 Å, the larger values being associated with the y parameters. Since the heavier atoms have the larger temperature factors, the uncertainties are roughly the same for all five heavy atoms. It seems appropriate, then, to estimate the limit of error as 0.0035 Å; this corresponds to limits of error of about 0.005 Å in the apparent bond distances and 0.3° in the bond angles. The estimated limit of error in the apparent hydrogen-atom positions is 0.06 Å.

The limits of error in the temperature-factor parameters, judging from the appearances of the difference maps, are probably around 0.1 \AA^2 in B units.

Discussion of the results

(i) Bond distances and angles

The bond distances and angles involving the heavy atoms are listed in Table 7 and shown in Fig. 4; the values in parentheses have been corrected for the rocking effects implied from the temperature-factor

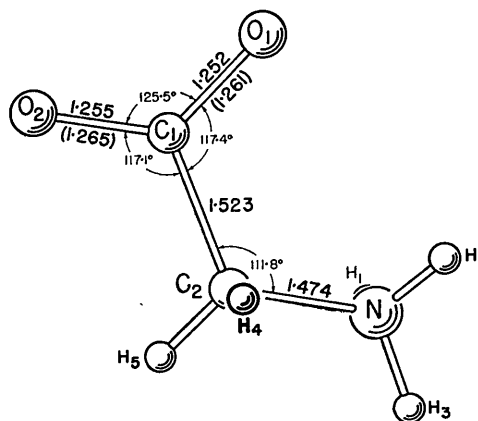


Fig. 4. Dimensions of the glycine molecule. The values in parentheses have been corrected for libration effects implied by the temperature-factor anisotropies.

Table 7. *Bond distances and bond angles*

Bond	This investigation	Albrecht & Corey	Angle	This investigation	Albrecht & Corey
C ₁ -O ₁	1.252 (1.261) Å	1.27 Å	C ₂ -C ₁ -O ₁	117.4°	119°
C ₁ -O ₂	1.255 (1.265)	1.25	C ₂ -C ₁ -O ₂	117.1	119
C ₁ -C ₂	1.523	1.52	O ₁ -C ₁ -O ₂	125.5	122
C ₂ -N	1.474	1.39	C ₁ -C ₂ -N	111.8	112

anisotropics.* For comparison, the values reported by Albrecht & Corey are also listed in Table 7. With the exception of the C-N bond, the molecular dimensions found in the present investigation are close to those reported by Albrecht & Corey; the C-N bond, on the other hand, is longer by 0.08 Å and is close to the values reported for other amino acids.

Bond distances and angles involving the hydrogen atoms are listed in Table 8. Although all the bond

Table 8. *Bond distances and angles involving the hydrogen atoms*

N-H ₁	0.92 Å	C ₂ -N-H ₁	110°
N-H ₂	0.85	C ₂ -N-H ₂	110
N-H ₃	0.85	C ₂ -N-H ₃	111
		H ₁ -N-H ₂	106
		H ₁ -N-H ₃	106
		H ₂ -N-H ₃	113
C ₂ -H ₄	0.91	C ₁ -C ₂ -H ₄	106
C ₂ -H ₅	0.91	C ₁ -C ₂ -H ₅	108
		N-C ₂ -H ₄	109
		N-C ₂ -H ₅	110
		H ₄ -C ₂ -H ₅	112

angles involving hydrogen atoms are within 4° of the expected tetrahedral values, the average C-H distance—0.91 Å—and the average N-H distance—0.87 Å—are both 0.16 Å shorter than the normal single-bond distances. This difference is far greater than the estimated limit of error in the present determination (0.06 Å), and seems to indicate that the effective centers of electron density of the hydrogen atoms are significantly displaced toward the adjacent atoms. Similar effects were observed in a two-dimensional analysis of salicylic acid (Cochran, 1953); on the other hand, a two-dimensional analysis of α -pyridone (Penfold, 1953) indicated more normal C-H and N-H distances, in agreement with the later findings of McDonald (1956) and Cochran (1956).

The electron densities associated with the hydrogen atoms in glycine (see Fig. 2) are roughly spherical; furthermore, the density distributions along the N-H

* The directions of maximum temperature factor for the two oxygen atoms lie within 5° of normal to the respective C₁-O bonds, and it has accordingly seemed appropriate to subtract the total temperature factor of C₁ from those of the oxygen atoms and to assume that the remainders represent the rocking motions of the oxygen atoms relative to C₁. For each oxygen atom the principal component of this remainder has a magnitude of approximately 2 Å² in *B* units and a direction approximately normal to the plane of the carboxyl group; only this principal component has been taken into account in calculating the bond shortening. No corrections were applied to the C₁-C₂ and C₂-N distances.

and C-H bonds are approximately Gaussian. It is apparent that the negative region around the nitrogen atom in the difference map, caused by the failure to compensate for the charge effect, could alter only slightly the apparent positions of H₁, H₂, and H₃; no such effect pertains to H₄ and H₅. It is perhaps worth while to point out that in the present case, because the structure has a relatively small over-all temperature factor, it has been possible to obtain intensity data out to a relatively small spacing and the electron densities associated with the hydrogen atoms have been clearly separated from the temperature-factor effects of the heavy atoms; such a separation is often not possible with copper intensity data. Indeed, if the apparent shortening of the C-H and N-H bonds is anomalous, owing to incorrect collection or interpretation of the data, the reason for the anomaly is not readily apparent.

Atoms C₁, C₂, O₁ and O₂ are nearly coplanar, the central carbon atom C₁ being only 0.005 Å from the least-squares plane (calculated by the method of Schomaker, Waser, Marsh & Bergman (1958) with all four atoms weighted equally). The nitrogen atom is 0.436 Å out of this plane.

(ii) *The hydrogen bonding*

Hydrogen-bond distances and angles are given in Table 9. The hydrogen-bond arrangement is essentially

Table 9. *Hydrogen-bond distances and angles*

N-O ₁ '	2.768 Å	C ₂ -N-O ₁ '	118.0°
H ₁ -O ₁	1.87	N-H ₁ -O ₁	168
N-O ₂ '	2.850	C ₂ -N-O ₂ '	116.1
H ₂ -O ₂ '	2.03	N-H ₂ -O ₂ '	163
N-O ₁ ''	2.949	C ₂ -N-O ₁ ''	155.8
N-O ₂ ''	3.074	C ₂ -N-O ₂ ''	93.0
H ₃ -O ₁ ''	2.44	N-H ₃ -O ₁ ''	120
H ₃ -O ₂ ''	2.29	N-H ₃ -O ₂ ''	155

that described by Albrecht & Corey. Strong N-H...O bonds involving H₁ and H₂ hold the molecules together in the *ac* plane (see Fig. 5); the third hydrogen atom of the ammonium group, H₃, is situated between two neighboring oxygen atoms, forming a weak, 'bifurcated' bond in the *b* direction. The geometry of the 'bifurcated' bond is essentially that pictured by Albrecht & Corey in their Fig. 13: although the nitrogen atom is closer to the neighboring O₁ atom than to the O₂ atom, the angle of attack is less favorable; accordingly, the hydrogen atom H₃ is appreciably closer to O₂ than to O₁.

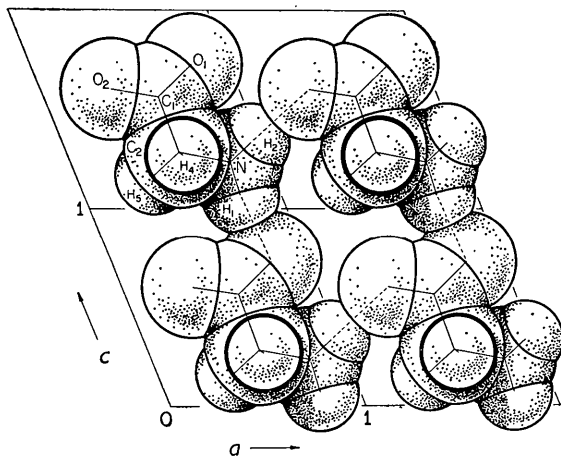


Fig. 5. A packing drawing showing the hydrogen bonds that link the molecules into layers parallel to (010).

(iii) *The temperature factors*

The temperature-factor parameters for the five heavy atoms were transformed into parameters related to the principal axes of the ellipsoids of vibration (Rollett & Davies, 1955); the magnitudes of these principal axes and their direction cosines relative to the unit-cell axes are listed in Table 10. Schematic

Table 10. *Magnitudes and directions of the principal axes of the vibration ellipsoids*

Atom	Axis i	B_i	g_1^i	g_2^i	g_3^i
C_1	1	1.564	0.833	-0.417	0.647
	2	1.200	0.426	0.903	0.213
	3	0.884	0.354	-0.105	-0.732
C_2	1	1.960	0.401	0.899	0.311
	2	1.320	0.876	-0.433	0.523
	3	1.052	0.268	0.067	-0.794
N	1	2.352	0.406	0.914	0.176
	2	1.324	0.889	-0.401	0.534
	3	0.964	0.213	-0.064	-0.827
O_1	1	3.084	0.378	0.923	0.206
	2	1.644	0.920	-0.366	0.209
	3	1.072	-0.104	0.119	-0.955
O_2	1	3.260	-0.002	0.984	-0.163
	2	2.136	0.751	0.118	0.882
	3	1.056	0.661	-0.130	-0.442

drawings of the ellipsoids viewed down the b and a axes are shown in Fig. 6.

The vibrations implied by the temperature-factor parameters are in general agreement with those that would be expected on the basis of the geometry of the molecule and the intermolecular packing. The smallest component for each atom is in a direction approximately parallel to the main axis of the molecule and to the shortest N-H...O hydrogen bond; the largest component of vibration for each atom except C_1 is approximately perpendicular to the plane of the

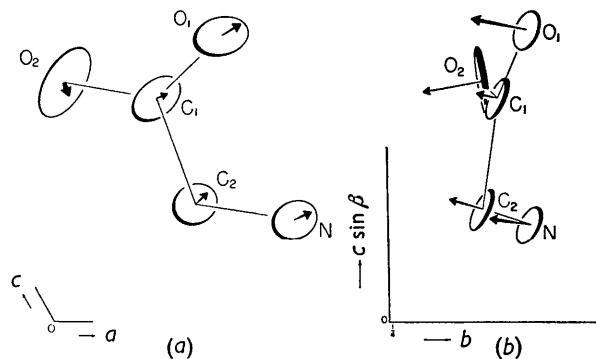


Fig. 6. (a) Schematic drawing of the ellipsoids of thermal vibration viewed along the b axis. The arrows are drawn in the directions of the principal axes most nearly perpendicular to the molecular plane; for all atoms except C_1 , they correspond to the major axes. The lengths of the arrows are proportional to the magnitudes of the corresponding axes projected on to the plane of the drawing. (b) Schematic drawing of the ellipsoids of thermal vibration viewed along the a axis.

molecule—that is, perpendicular to the plane of the strong hydrogen bonding. The central carbon atom C_1 , which is covalently bonded to three heavy atoms, has the smallest average vibration, followed by C_2 and the nitrogen atom; the oxygen atoms, which are each bonded (covalently) to only one atom, have the largest vibrations. Furthermore, O_2 , which is involved in the longer and hence weaker of the two strong hydrogen bonds, has a larger vibration than O_1 .

The relatively large anisotropy of C_1 in the ac plane is surprising; it is conceivable that it is an artifact caused by the electron density associated with the C_1-O_1 bond, although the use of extensive high-angle data has presumably minimized any such effect.

I should like to express my gratitude to Prof. Robert B. Corey for his continuing interest and encouragement throughout the course of this investigation.

References

- ALBRECHT, G. & COREY, R. B. (1939). *J. Amer. Chem. Soc.* **61**, 1087.
 COCHRAN, W. (1953). *Acta Cryst.* **6**, 260.
 COCHRAN, W. (1956). *Acta Cryst.* **9**, 924.
 CURTIUS, T. (1882). *J. Prakt. Chem.* **26**, 158.
 HOERNI, J. A. & IBERS, J. A. (1954). *Acta Cryst.* **7**, 744.
 LEUNG, Y., MARSH, R. E. & SCHOMAKER, V. (1957). *Acta Cryst.* **10**, 650.
 LOW, B. W. & RICHARDS, F. M. (1952). *J. Amer. Chem. Soc.* **74**, 1660.
 McDONALD, T. R. R. (1956). *Acta Cryst.* **9**, 162.
 PASTERNAK, R. A. (1956). *Acta Cryst.* **9**, 341.
 PENFOLD, B. R. (1953). *Acta Cryst.* **6**, 591.
 ROLLETT, J. S. & DAVIES, D. R. (1955). *Acta Cryst.* **8**, 125.
 SCHOMAKER, V., WASER, J., MARSH, R. E. & BERGMAN, B. G. (1958). *Acta Cryst.* (in the press).